

Title	Superstring Theory and Triple Systems (Algebra and Computer Science)
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Citation	数理解析研究所講究録 (2014), 1873: 114-121
Issue Date	2014-01
URL	http://hdl.handle.net/2433/195507
Right	
Type	Departmental Bulletin Paper
Textversion	publisher

Superstring Theory and Triple Systems

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1 Introduction

It has been expected that there exists M-theory, which unifies string theories. In M-theory, some structures of 3-algebras were found recently. First, it was found that by using $u(N) \oplus u(N)$ Hermitian 3-algebra, we can describe a low energy effective action of N coincident supermembranes [1–5], which are fundamental objects in M-theory.

Second, recent studies have indicated that there also exist structures of 3-algebras in the Green-Schwartz supermembrane action, which defines full perturbative dynamics of a supermembrane. It had not been clear whether the total supermembrane action including fermions has structures of 3-algebras, whereas the bosonic part of the action can be described by using a tri-linear bracket, called Nambu bracket [6, 7], which is a generalization of Poisson bracket. If we fix to a light-cone gauge, the total action can be described by using Poisson bracket, that is, only structures of Lie algebra are left in this gauge [8]. However, it was shown under an approximation that the total action can be described by Nambu bracket if we fix to a semi-light-cone gauge [9]. In this gauge, the eleven dimensional space-time of M-theory is manifest in the supermembrane action, whereas only ten dimensional part is manifest in the light-cone gauge. Moreover, 3-algebra models of M-theory itself were proposed and have been studied in [9–13].

The hermitian (ϵ, δ) -Freudenthal-Kantor triple systems [14–36] are generalizations of the hermitian 3-algebras [1–9, 37–71]. The hermitian 3-algebras are special cases, where $K(a, b) = 0$ or equivalently, $\langle abc \rangle = -\langle cba \rangle$, of the hermitian $(-1, -1)$ -Freudenthal-Kantor triple systems of second order. And the hermitian 3-algebras are classified into the $u(N) \oplus u(M)$ and $sp(2N) \oplus u(1)$ hermitian 3-algebras [13, 43, 45, 46, 52]. Therefore, it is natural to extend these triple systems to more general hermitian $(-1, -1)$ -Freudenthal-Kantor triple systems or hermitian generalized Jordan triple systems.

In the following section, we summarize some results concerning with the generalization of the hermitian 3-algebras in M-theory [72, 73].

2 Definitions

Let us start with a definition of a \ast -(ϵ, δ)-Freudenthal-Kantor triple system.

Definition. A triple system U is said to be a \ast -(ϵ, δ)-Freudenthal-Kantor triple system if relations (0)–(iv) satisfy;

- 0) U is a Banach space,

- i) $[L(a, b), L(c, d)] = L(\langle abc \rangle, d) + \varepsilon L(c, \langle bad \rangle)$,
 ii) $K(\langle abc \rangle, d) + K(c, \langle abd \rangle) + \delta K(a, K(c, d)b) = 0$,
 where $L(a, b)c = \langle abc \rangle$ and $K(a, b)c = \langle acb \rangle - \delta \langle bca \rangle$, $\varepsilon = \pm 1, \delta = \pm 1$,
 iii) $\langle xyz \rangle$ is \mathbf{C} -linear operator on x, z and \mathbf{C} -anti-linear operator on y ,
 iv) $\langle abc \rangle$ continuous with respect to a norm $\| \cdot \|$ that is, there exists $K > 0$ such that

$$\| \langle xxx \rangle \| \leq K \|x\|^3 \text{ for all } x \in U.$$

Furthermore, a \ast -(ϵ, δ)-Freudenthal-Kantor triple system is said to be hermitian if it satisfies the following condition,

- v) all operator $L(x, y)$ is a positive hermitian operator with a hermitian metric

$$(x, y) = \text{tr } L(x, y),$$

that is, $(L(x, y)z, w) = (z, L^*(x, y)w)$, and $(x, y) = \overline{(y, x)}$.

Let U be a \ast -(ϵ, δ)-Freudenthal-Kantor triple system. Then we may define the notation of tripotent as follows.

Definition. It is said to be a tripotent of U if

$$\langle ccc \rangle = c, \quad c \in U.$$

3 Tripotent basis

In this section, we give decomposition theorems based on the tripotent basis.

Theorem 1.1. Let U be a hermitian $(-1, \delta)$ -Freudenthal-Kantor triple system. If $W \subset U$ is flat (that is, $L(x, y) = L(y, x)$ for all $x, y \in W$), then we have a decomposition,

$$W = \mathbf{R}e_1 \oplus \cdots \oplus \mathbf{R}e_n$$

where e_i are tripotents or bitripotents.

Proof. See [72, 73].

We define the odd power of x inductively as follows;

$$\begin{aligned} x^{(3)} &:= \langle xxx \rangle, \\ x^{(2n+1)} &:= \langle xx^{(2n-1)}x \rangle. \end{aligned}$$

By using this theorem, we have

Theorem 1.2. Let U be a hermitian $(-1, \delta)$ -Freudenthal-Kantor triple system. Then every $x \in U$ can be written uniquely

$$x = \lambda_1 e_1 + \lambda_2 e_2 + \cdots + \lambda_n e_n,$$

where the e_i are tripotents or bitripotents, which are linear combinations of power of x , and the λ_i satisfy

$$0 < \lambda_1 \leq \lambda_2 \leq \cdots \leq \lambda_n.$$

Poof. See [72, 73].

4 Peirce decomposition

In this section, we give a theorem on a Peirce decomposition of a $\ast(-1, -1)$ -Freudenthal-Kantor triple system equipped with the tripotent $\langle ccc \rangle = c$.

Theorem 2.1. Let U be a $\ast(-1, -1)$ -Freudenthal-Kantor triple system. Then, we have a decomposition with respect to a tripotent c (i.e., $\langle ccc \rangle = c$) as follows

$$U = U_1(c) \oplus U_{\frac{1}{2}}(c) \oplus U_0(c),$$

where

$$U_1(c) = \{x | (L(c, c) + R(c, c))x = 0, (R(c, c) - Id)x \neq 0\},$$

$$U_{\frac{1}{2}}(c) = \{x | (L(c, c) + R(c, c))x \neq 0, (R(c, c) - Id)x = 0\},$$

$$U_0(c) = \{x | (L(c, c) + R(c, c))x = 0, (R(c, c) - Id)x = 0\}.$$

Poof. See [72, 73].

5 Generalized hermitian 3-algebra

In this section, we extend the $u(N) \oplus u(M)$ 3-algebras to a hermitian $(-1, -1)$ -Freudenthal-Kantor triple system.

Let $D_{N,M}^*$ be the set of all $N \times M$ matrices with operation

$$\langle xyz \rangle = x\bar{y}^T z - z\bar{y}^T x + zx^T \bar{y},$$

where x^T and \bar{x} mean transpose and conjugation of x , respectively.

Then $D_{N,M}^*$ is a hermitian $(-1, -1)$ -Freudenthal-Kantor triple system. In fact, it satisfies the conditions (0), (i), (ii), (iii), (iv) and (v). This is an extension of the $u(N) \oplus u(M)$ hermitian 3-algebra, $\langle xyz \rangle = x\bar{y}^T z - z\bar{y}^T x$, which is a basis for the effective action of the multiple membranes in M-theory.

One of the tripotents is given by

$$c = \begin{pmatrix} Id & 0 \\ 0 & 0 \end{pmatrix},$$

where Id is a $n \times n$ identity matrix ($n \leq N, M$). Because any element is decomposed as

$$x = \begin{pmatrix} A & B \\ C & D \end{pmatrix} = \begin{pmatrix} \frac{1}{2}(A - A^T) & B \\ 0 & D \end{pmatrix} + \begin{pmatrix} \frac{1}{2}(A + A^T) & 0 \\ C & 0 \end{pmatrix},$$

the Peirce decomposition is given by

$$U_1(c) = \left\{ \begin{pmatrix} \frac{1}{2}(A - A^T) & B \\ 0 & D \end{pmatrix} \right\},$$

$$U_{\frac{1}{2}}(c) = \left\{ \begin{pmatrix} \frac{1}{2}(A + A^T) & 0 \\ C & 0 \end{pmatrix} \right\},$$

$$U_0(c) = 0.$$

As in Theorem 1.1, we can expand any element as $x = \sum(\lambda_{ij} E_{ij} + \mu_{ij} \sqrt{-1} E_{ij})$, where E_{ij} means that (i, j) element is 1 and other element is zero, and E_{ij} and $\sqrt{-1} E_{ij}$ are tripotents, i.e., $\langle E_{ij} E_{ij} E_{ij} \rangle = E_{ij}$, and $\langle \sqrt{-1} E_{ij} \sqrt{-1} E_{ij} \sqrt{-1} E_{ij} \rangle = \sqrt{-1} E_{ij}$.

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